

Infection propagator approach to compute epidemic thresholds on temporal networks: impact of immunity and of limited temporal resolution[★]

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Abstract. The epidemic threshold of a spreading process indicates the condition for the occurrence of the wide spreading regime, thus representing a predictor of the network vulnerability to the epidemic. Such threshold depends on the natural history of the disease and on the pattern of contacts of the network with its time variation. Based on the theoretical framework introduced in [E. Valdano, L. Ferreri, C. Poletto, V. Colizza, Phys. Rev. X **5**, 21005 (2015)] for a susceptible-infectious-susceptible model, we formulate here an infection propagator approach to compute the epidemic threshold accounting for more realistic effects regarding a varying force of infection per contact, the presence of immunity, and a limited time resolution of the temporal network. We apply the approach to two temporal network models and an empirical dataset of school contacts. We find that permanent or temporary immunity do not affect the estimation of the epidemic threshold through the infection propagator approach. Comparisons with numerical results show the good agreement of the analytical predictions. Aggregating the temporal network rapidly deteriorates the predictions, except for slow diseases once the heterogeneity of the links is preserved. Weight-topology correlations are found to be the critical factor to be preserved to improve accuracy in the prediction.

1 Introduction

The concept of epidemic threshold is fundamental in infectious disease modeling [1,2]. When a pathogen is seeded in a population, a critical transmissibility exists below which the spread rapidly ceases. Such a threshold is a combined property of the disease natural history and of the network of interactions along which transmission can occur. In the physics literature such interplay has been typically studied for the family of susceptible-infected-susceptible and susceptible-infected-recovered models on networks [3–11]. Several analytical approaches based, for instance, on the heterogeneous mean field approximation [3], on percolation theory [6,7] and on Markov processes [4,5] have been developed to study the transition from early extinction to epidemic.

Extensive work has been done under the assumption of spreading time scales either much slower or much faster than the one characteristic of the underlying network –

the two regimes called annealed and quenched, respectively [3,5]. In recent years, the massive amount of empirical information on networks has showed that such assumption does not hold in many cases [12–19] and that the network dynamics presents features (e.g. memory, bursty activation, heterogeneities in node activity) affecting the resulting spreading processes [13,14,18,20–28].

The majority of studies addressing so far the impact of network dynamics on the epidemic spread through the analytical calculation of the epidemic threshold are all based on synthetic models of the network evolution, valid under context-specific assumptions [21,24–29]. To fill this gap, we have introduced in reference [30] a method to compute the epidemic threshold for a susceptible-infectious-susceptible (SIS) process on a generic discrete-time temporal network, assuming the knowledge of its sequence of adjacency matrices. The approach is rooted in a multi-layer representation [31,32] of the temporal network that preserves the network causality. It employs a tensor formulation that integrates both spreading and network dynamics and allows for the analytical solution of the linearized Markov chain description of the spreading process. Such framework extends in this way the quenched approach to the time-varying case, through a multilayer transformation.

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The lack of assumptions on the network substrate makes such a tool a candidate for assessing the vulnerability to epidemic invasion of real systems for which time-varying contact data relevant for epidemic transmissions are collected [13–19,33–36]. At the same time, it allows a systematic exploration of the structural and temporal factors characterizing the time-evolving network that are responsible for sustained spreading. To allow the use of this framework to a variety of different settings and epidemic conditions, we assess here the applicability of the approach in describing realistic diseases and its robustness with respect to network properties induced by data collection procedures and availability. By considering empirical and synthetic model contact data, we discuss how a varying force of infection along a given link and its direction impact the computation of the threshold for a SIS dynamics. Then, we formulate the approach for more realistic disease natural histories, considering susceptible-infectious-recovered (SIR) and susceptible-infectious-recovered-susceptible (SIRS) compartmental models. This allows us to account for an additional ingredient – immunity following infection, either permanent or temporary – representing an important feature of many diseases. Finally, we address the problem of limited temporal resolution in the knowledge or availability of the network dynamics, for which contacts occurring within a given time interval are aggregated [37]. By focusing on an empirical network of time-varying contacts among individuals at school, we quantify the accuracy and reliability of the estimation of the epidemic threshold testing increasing aggregations, to provide quantitative and qualitative information on the specific temporal structures responsible for observed biases.

2 Infection propagator approach for a weighted directed temporal network

We consider a SIS model [1,2] where hosts, represented by nodes in the network, can be either in the susceptible or infectious state. We assume the process to unfold in discrete time on a weighted directed temporal network, comprising a finite number T of snapshots, each one with a weighted adjacency matrix \mathbf{W}_t . The entry $W_{t,ij}$ encodes the weight of the directed link from i to j at time step t . At each time step, infectious nodes spontaneously recover with probability μ , returning to the susceptible state. While infectious, nodes can transmit the infection to susceptible neighbors with a probability that depends both on the weight of the link and on the intrinsic transmissibility of the pathogen λ , representing the probability of transmission when link weight is equal to 1. We model this by introducing a transmission matrix $\mathbf{\Lambda}_t$, function of both λ and \mathbf{W}_t , that encodes transmission probabilities. When the network is unweighted ($W_{t,ij} = A_{t,ij} = 0, 1$), the entries of the matrix $\mathbf{\Lambda}_t$ are simply given by $\Lambda_{t,ij} = \lambda A_{t,ij}$. If the network is weighted, several choices are possible to model transmission along the weighted link. Here we

consider a binomial process for the infection

$$A_{t,ij} = 1 - (1 - \lambda)^{W_{t,ij}}, \quad (1)$$

as it is typically assumed, for example, in the spread of livestock infections between premises where the weight represents the number of animals moved between farms [38], or in the context of social networks, where weight may encode force-of-infection in terms of either contact frequency or duration [39].

We start from the microscopic Markov chain approach, or *quenched mean field approach*, developed for static networks [4,5]. According to this, the equations describing the SIS propagation on a generic static network with N nodes and adjacency matrix \mathbf{W} are

$$p_{t,i} = 1 - [1 - (1 - \mu) p_{t-1,i}] \prod_j (1 - \Lambda_{ji} p_{t-1,j}), \quad (2)$$

where $p_{t,i}$ is the probability that node i is in the infectious state at time t , and $\mathbf{\Lambda}$ is the static transmission matrix. We remark that equation (2) relies on the assumption that no dynamical correlations exist among infection probabilities of neighbouring nodes [40].

The microscopic Markov chain model of equation (2) is widely adopted in different fields [8,41]. For both directed and undirected networks [42,43] the study of its asymptotic state yields the derivation of the epidemic threshold in terms of the spectral radius of the transmission matrix $\rho[\mathbf{\Lambda}]$, namely the modulus of the largest eigenvalue of $\mathbf{\Lambda}$. The threshold is the value of λ for which the following holds: $\rho[\mathbf{\Lambda}] = \mu$. In the unweighted case, this is equivalent to the well-known relation $(\lambda/\mu) = 1/\rho[\mathbf{A}]$, where $\rho[\mathbf{A}]$ is the spectral radius of the adjacency matrix [4,5].

In order to extend this approach to temporal networks we need to take into account the time dependence of $\mathbf{\Lambda}$. The Markov chain equations of the process read in this case:

$$p_{t,i} = 1 - [1 - (1 - \mu) p_{t-1,i}] \prod_j (1 - \Lambda_{t-1,ji} p_{t-1,j}). \quad (3)$$

We enforce the existence of the asymptotic solution of the infection process in a generic temporal network by imposing periodic boundary conditions for network dynamics, i.e. $\mathbf{W}_{T+1} \equiv \mathbf{W}_1$. Given that T is arbitrary, this causes no loss in generality. We also tested that it would affect the epidemic threshold estimation only for rather small values of T , also when complex temporal dynamics are considered [30]. Contrary to the static case, now the asymptotic solutions of equation (3) are periodic of period T .

We develop the formalism introduced in reference [30], and compute the epidemic threshold for a generic weighted directed temporal network. We define a new representation of the SIS dynamics on a temporal network by employing a multi-layer representation [31,32,44]. We map the temporal network to the tensor space $\mathbb{R}^N \otimes \mathbb{R}^T$, where each node is identified by the pair of indices (i, t) , corresponding to the node label i and the time frame t , respectively. Layer t thus contains the images of the nodes

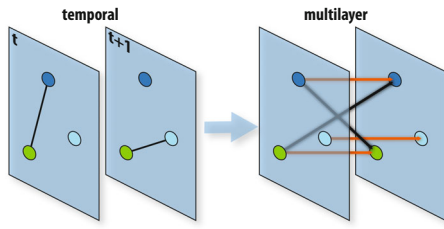


Fig. 1. Multi-layer representation of the temporal network. We consider a temporal network of 3 nodes and 2 time steps (left). This network is mapped onto a multi-layer representation (right). Each node points to its future self (orange arrows), and to the future image of its present neighbors (black arrows).

corresponding to time step t . The specific multi-layer representation of the temporal network is built according to the following rules [30]:

- (i) each node, at time t , is connected to its future self-image at time $t + 1$;
- (ii) if i is connected to j at time t with weight w , then we connect i at time t to j at time $t + 1$, and j at time t to i at time $t + 1$, both with weight w .

These rules define a tensor representation of a weighted multilayer network [32,45]. We stress that no links connect nodes on the same layer, as contacts in the temporal networks are mapped onto the inter-layer links of the multilayer object (rule (ii)). The resulting network is thus multipartite, since only pairs of nodes belonging to different layers are linked together. Figure 1 provides a schematic representation of this multilayer mapping. The adjacency representation of the resulting multilayer network has as entries $\hat{W}_{tt',ij} = \delta_{t,t'+1} [\delta_{ij} + W_{t,ij}]$. The proposed mapping from the network temporal sequence to a multilayer object provides an ad hoc representation of the temporal network that preserves the causality of the temporal network and that it lends itself to the integration of the infection and recovery processes. The transformation for the links (rule (ii)) is similar to the one introduced in reference [46], and is here introduced to model the infection process along a time-stamped link. In addition, we also need to consider the connection between each node and its future self (rule (i)) to model the recovery process of each infected node. We can therefore define the transmission tensor \mathbf{M} , whose entries are defined as:

$$M_{tt',ij} = \delta_{t,t'+1} [(1 - \mu) \delta_{ij} + A_{t,ij}]. \quad (4)$$

\mathbf{M} contains the transmission terms \mathbf{A}_t and the recovery term. \mathbf{M} also introduces a simplified expression for equation (3). Using the supra-adjacency matrix formalism [31,47,48], we can flatten out the multilayer representation using the following mapping: $(i, t) \rightarrow \alpha = Nt + i$, with α running in $\{1, \dots, NT\}$, allowing us to write \mathbf{M} in matrix form

$$\mathbf{M} = \begin{pmatrix} 0 & 1-\mu+\mathbf{A}_1 & 0 & \dots & 0 \\ 0 & 0 & 1-\mu+\mathbf{A}_2 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & 1-\mu+\mathbf{A}_{T-1} \\ 1-\mu+\mathbf{A}_T & 0 & 0 & \dots & 0 \end{pmatrix}.$$

\mathbf{M} provides a representation of the topological and temporal dimensions underlying the dynamics of equation (3), in terms of a $NT \times NT$ transmission matrix that encodes both pathogen transmission and recovery. The Markov process is now represented in \mathbb{R}^{NT} by the state vector $\hat{p}_\alpha(\tau)$, i.e. the probability of each node to be infectious at each time step t included in a 1-period long interval, $[\tau T, (\tau + 1)T]$. Consistently, equation (3) becomes

$$\hat{p}_\alpha(\tau) = 1 - \prod_{\beta} [1 - M_{\beta\alpha} \hat{p}_\beta(\tau - 1)]. \quad (5)$$

Given that vector \hat{p} encodes a 1-period configuration, the T -periodic asymptotic state of the SIS process is now mapped into the steady state $\hat{p}_\alpha(\tau) = \hat{p}_\alpha(\tau - 1)$. The latter can be recovered as solution of the equilibrium equation:

$$\hat{p}_\alpha = 1 - \prod_{\beta} (1 - M_{\beta\alpha} \hat{p}_\beta), \quad (6)$$

that is formally the same as the stationary condition imposed on equation (2) for the static network case and is similar to Markov chain approaches used to solve contagion processes in multiplex and interconnected networks [47–49]. Given that equation (6) formally describes a diffusion process on a static network of NT nodes, we can then follow [4,5] and linearize equation (6) recovering the necessary and sufficient condition for the asymptotically stable disease-free solution, $\rho[\mathbf{M}] < 1$ [50]. This yields the threshold condition

$$\rho[\mathbf{M}] = 1 \quad (7)$$

for the critical value of λ above which the transmission becomes epidemic [4,5,47–49]. Given the block structure of the matrix, it is possible to simplify the computation of the spectral radius of \mathbf{M} [51]:

$$\rho[\mathbf{M}] = \rho[\mathbf{P}]^{1/T} \quad (8)$$

where

$$\mathbf{P} = \prod_{t=1}^T (1 - \mu + \mathbf{A}_t). \quad (9)$$

In the case of unweighted undirected network, \mathbf{P} becomes $\mathbf{P} = \prod_{t=1}^T (1 - \mu + \lambda \mathbf{A}_t)$ [30]. This matrix has an important physical interpretation. Let us consider a time-respecting path from i to j , lasting T time steps and containing a jumps and $T - a$ waiting times. We associate to this path the weight $\lambda^a (1 - \mu)^{T-a}$, representing the probability that the infection propagates along that path, from i infectious at time $t = 1$ to j infectious at time $t = T$. The entry P_{ij} is then the sum of all the time-respecting paths going from $(i, t = 1)$ to $(j, t = T)$, each weighted as described. Therefore, it represents the total probability of j being infectious at time $t = T$, given that the infection originated in i infectious at time $t = 1$. This is valid in the limit of small probabilities and non-interaction among paths. \mathbf{P} thus describes the infection propagation around the disease free state (i.e. $p \simeq 0$) and within the quenched

mean field framework where interactions among paths are disregarded. In light of this interpretation, we call \mathbf{P} *infection propagator*. The accessibility matrix, defined in reference [52], is a particular case of infection propagator, when $\lambda = \mu = 1$, i.e., when the spreading process is a deterministic exploration of the temporal network. The generalization to weighted networks is straightforward once the force of transmission on each link of equation (1) is taken into account.

3 Infection propagator approach for SIRS and SIR dynamics

Many pathogens leave recovered individuals immune to reinfection. Such immunity may last indefinitely, or for a limited amount of time and is modeled through the introduction of an additional compartment, the recovered (R) state [1,2]. Infectious nodes enter the recovered state with probability μ , becoming immune to re-infection. We also consider that they leave this state with probability ω , returning to the susceptible state. Any value of $\omega > 0$ describes a SIRS model, characterized by an average immunity period $1/\omega$. $\omega = 0$ corresponds instead to the SIR model, where immunity is assumed to be permanent. Markov chain equations for the SIRS model are as follows:

$$\begin{cases} p_{t,i} = (1 - \mu) p_{t-1,i} \\ \quad + (1 - p_{t-1,i} - q_{t-1,i}) \left[1 - \prod_j (1 - \Lambda_{t-1,ji} p_{t-1,j}) \right]; \\ q_{t,i} = \mu p_{t-1,i} + (1 - \omega) q_{t-1,i}. \end{cases} \quad (10)$$

In addition to $p_{t,i}$, we define $q_{t,i}$ as the probability of being in the recovered state at time t . The computation of the threshold is equivalent to the study of the stability of the disease-free state $p_{t,i} = 0$ [4,5]. Equations are therefore linearized around that point, making all quadratic terms disappear. In the case of the SIRS model, the disease-free state is $p_{t,i} = q_{t,i} = 0$ whose stability is now studied by linearizing in both $p_{t,i}$ and $q_{t,i}$. The linearized form of equation (10) is the following:

$$\begin{cases} p_{t,i} \approx \sum_j (\Lambda_{t-1,ji} + (1 - \mu) \delta_{ij}) p_{t-1,j} \\ q_{t,i} = \mu p_{t-1,i} + (1 - \omega) q_{t-1,i}. \end{cases} \quad (11)$$

From this we see that the equation for \mathbf{p}_t no longer contains \mathbf{q}_t , showing that the evolution of the number of infected around the disease-free state does not depend on the recovered individuals. This in turns implies that the recovered compartment does not impact the epidemic threshold, as in the static case [53,54]. As a result, the same infection propagator describing the SIS dynamics (Eq. (9)) can be used to compute the epidemic threshold of a SIRS compartmental model on a time-varying network.

The SIR model can be considered as a limiting case of the SIRS dynamics ($\omega \rightarrow 0$). The infection propagator for the SIRS model does not depend on the probability of waning of immunity ω , as it only contains expressions in

terms of λ and μ . The threshold computed with the infection propagator approach for the SIRS therefore holds for any arbitrarily small ω . As a result, we can safely perform the following limit:

$$\lambda_{critical}^{SIR} = \lim_{\omega \rightarrow 0} \lambda_{critical}^{SIRS} = \lambda_{critical}^{SIRS} = \lambda_{critical}^{SIS}. \quad (12)$$

Both SIRS and SIR models thus have the same threshold as the SIS compartmental model, not being affected by the recovery compartment and the duration of the immunity period.

4 Application to empirical and synthetic model data

We test the validity and accuracy of our predictions by comparing them with the results of explicit microscopic stochastic numerical simulations of the SIR and SIRS processes. We consider two temporal network models and one empirical time-varying network. In the following subsections we describe the data and methods considered and the corresponding results.

4.1 Empirical and synthetic model data

We test our approach on two network models: ACTIVITY and BURSTY models. ACTIVITY is built from the activity-driven model proposed by Perra et al. in reference [25]. Each node is given an activity potential, drawn from a heterogeneous distribution. At each time step, nodes become active with a probability equal to their potentials. Active nodes establish m_{stub} (here $m_{stub} = 2$) connections with other nodes picked at random, and all links are renewed at every snapshot. We generate networks with $N = 1000$ nodes and $T = 20$ time snapshots, as in reference [30]. In addition, we explored size effects by considering networks with sizes ranging from $N = 10^2$ nodes to $N = 10^4$ nodes. Activity potentials are assigned through the relation $a = 1 - e^{-\eta x}$, ($\eta = 10$) and $x \sim x^{-\gamma}$ and $x \in [\epsilon, 1]$ ($\gamma = 2.8$ and $\epsilon = 3 \times 10^{-2}$). The obtained networks are characterized by a temporally uncorrelated sequence of snapshots yielding an aggregated network with heterogeneous topology.

BURSTY is obtained from the model introduced by Rocha and Blondel in reference [20]. Here, the probability of a node becoming active at time step t is sampled from the distribution $(t - t')^{-\alpha_1} e^{-\alpha_2(t-t')}$, where t' is the time that node was last active. We consider networks of size $N = 500$ and described by $T = 50$ time snapshots, generated with $\alpha_1 = 2$ and $\alpha_2 = 5 \times 10^{-4}$ as in reference [30]. The obtained networks account for a heterogeneous activation pattern describing a sequence of homogeneous networks where the inter-contact time is power-law distributed.

In addition to the synthetic models above, we consider an empirical time-evolving network constructed from records of face-to-face proximity interactions between individuals in a high school during one day, collected by

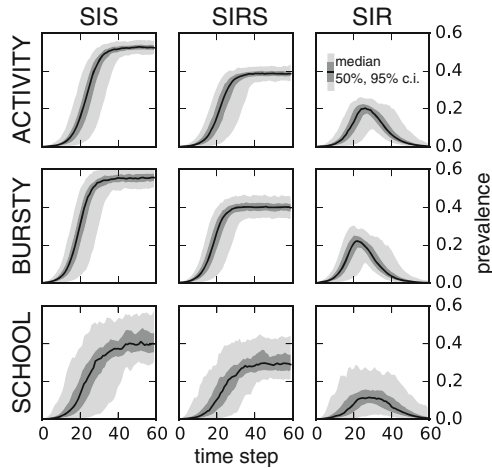


Fig. 2. Evolution of fraction of infected nodes (disease prevalence) for all considered disease models and networks. In all panels $\mu = 0.5$. For SIRS, $\omega = 0.75$. λ is set at $5/2$ of the epidemic threshold. This is an arbitrary value allowing us to show the typical behavior well above threshold. Median prevalence (solid black lines), and 50%, 95% confidence intervals (gray shaded areas) are computed on 200 runs, keeping only the ones not going extinct in the beginning.

Salathé et al. [55] (SCHOOL). This network comprises $N = 787$ nodes and we consider here $T = 42$ time snapshots, each one of 10 min. In Section 5 we will also examine the impact of a finer aggregation time.

4.2 Numerical simulations

We numerically simulate the disease diffusion of a SIR and of a SIRS infection dynamics on the above described networks. Simulations assume all individuals to be susceptible at the initial time, and are seeded with an infected node chosen at random on the network. At each time step, infectious nodes can transmit the disease with probability λ to their susceptible neighbors and recover with probability μ . Here we consider unweighted networks, for the sake of simplicity. Weighted networks will be addressed in the next section in the study of time aggregation of the evolving network. In the SIRS model, recovered nodes turn susceptible with probability ω . Results of the simulations are obtained after randomizing the initial seed and the time step of the T -sequence chosen as the initial time step, and they are obtained under the assumption of periodic boundary conditions for network evolution.

Figure 2 shows the typical evolution of an outbreak, for the three disease models considered (SIS, SIRS, and SIR) and for all networks. The above-threshold behavior is marked by a non-zero endemic state for both SIS and SIRS, whereas a SIR outbreak shows the usual bell-like trend, reaching always extinction at the end, as expected.

For the SIRS dynamics, following [56] we numerically identify the epidemic threshold as the value of the transmissibility λ for which the relative variation of the prevalence at equilibrium is maximal, as such variation would go

to infinity in the thermodynamic limit ($N \rightarrow \infty$), indicating a second order phase transition. We therefore measure the *variability* $\Delta = \sqrt{\langle i_T^2 \rangle - \langle i_T \rangle^2} / \langle i_T \rangle$ [54,56,57], where $i_T = \frac{1}{T} \sum_{t=1}^T i_{eq}(t)$ is the prevalence at equilibrium averaged over a period T . The endemic prevalence is computed using the quasistationary method [56,58]. We force the system to be in an active state; whenever it reaches the absorbing state with no infectious nodes, we sample one random configuration among the ones the system had visited while it was in the same snapshot, and restart the simulation from that configuration. After discarding an initial transient (3×10^3 iterations), we compute i_T for every period (for 5×10^5 iterations), and with those values we compute $\langle i_T \rangle$ and $\langle i_T^2 \rangle$. We then compare the value of λ corresponding to the peak of the variability Δ with the prediction for the epidemic threshold obtained from the infection propagator approach. The same method is used for the SIR dynamics, where the endemic prevalence is replaced by the final attack rate r , i.e. the fraction of nodes hit by the epidemic. We stress that variability has formally the same definition in both SIRS and SIR models, but it is computed using different observables. We are only interested in the position of the peak [56], as an indication of the epidemic threshold, and not its global behavior.

4.3 Threshold results

We consider a SIR dynamics on the three networks under study and explore two values of the infectious period, corresponding to $\mu = 0.2, 0.5$. For each μ , Figure 3 shows the behavior of the variability Δ normalized to its peak value Δ_{\max} as a function of the transmissibility λ . In all cases we find a very good agreement between our prediction (vertical dashed line) and the simulated epidemic threshold obtained from the peak position of Δ . The agreement is found for both network models, ACTIVITY and BURSTY, despite them being characterized by different topological and temporal heterogeneities, and for the empirical dataset SCHOOL. This last network features a more complex dynamics capturing the daily activities and interactions, with non-trivial temporal correlations and modular structures evolving in time [55]. Despite the approximations used to compute the epidemic threshold with the infection propagator approach, the results of Figure 3 indicate that the method is able to provide reliable and accurate predictions for the threshold behavior of systems characterized by different properties. We also note that the agreement is obtained independently of the values of the epidemic threshold: the threshold of the SCHOOL network is indeed approximately one order of magnitude smaller than the ones obtained in the two network models for the same SIR dynamics.

The choice of the number of nodes (N) and snapshots (T) in the models considered is arbitrary, but it does not impact our findings. In reference [30] we already studied the effect of T on the threshold, and observed that, after an initial transient, the epidemic threshold saturates around a stable value, showing that the optimal T has

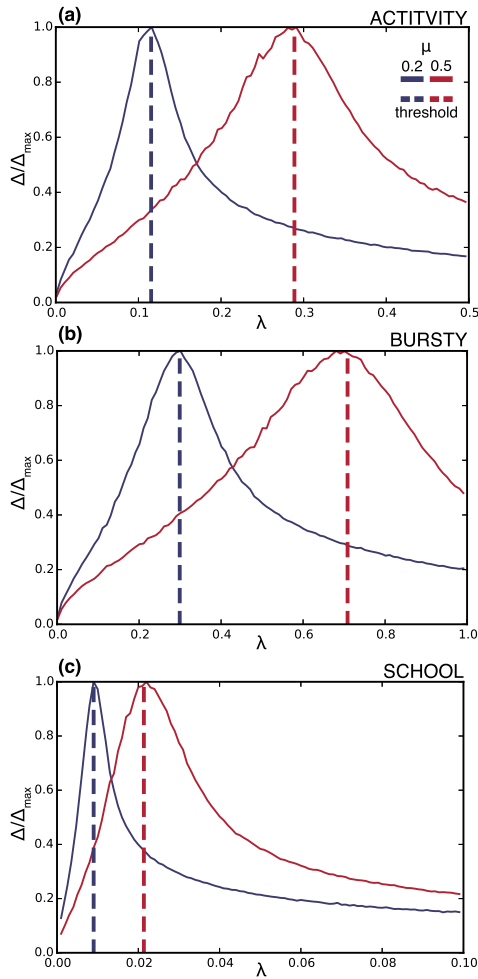


Fig. 3. SIR model. Comparison between the epidemic threshold and variability. The variability Δ of final attack rate r , normalized to its peak value (Δ/Δ_{\max}), is plotted against the transmissibility λ . We explore two different values of μ . Dashed vertical lines represent the threshold value for λ predicted by the infection propagator approach. (a) Shows the results for the ACTIVITY network, (b) for the BURSTY network, and (c) for the SCHOOL network.

been reached. Here, in Figure 4, we explore the effect of N for the ACTIVITY model, by building different instances of the model, for several values of N . We find that the median value of the threshold, computed over the instances, does not depend on the network size N . Fluctuations around the median, however, sharply decrease as N increases, as expected.

Results similar to those presented in Figure 3 are also obtained when considering a SIRS dynamics, characterized by the same values of the infectious period considered above and by three values of the probability of immunity waning ($\omega = 0.25, 0.5, 0.75$). For each temporal network, we numerically identify the value of the epidemic threshold as that corresponding to the peak of the normalized variability Δ/Δ_{\max} , and recover a good agreement with our analytical predictions, Figure 5. The addition of the transition from an immune state to a susceptible state

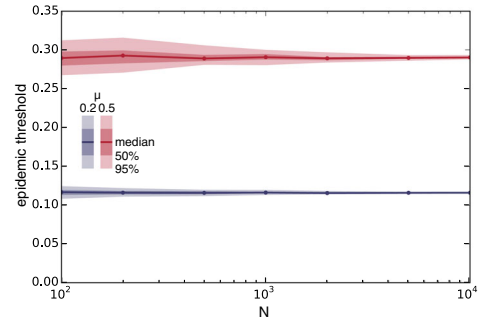


Fig. 4. Impact of the number of nodes on the epidemic threshold of the ACTIVITY model. We build ACTIVITY models for different values of N , and compute their epidemic threshold for $\mu = 0.2$. For each value of N (x -axis) we build 200 instances of the model, and compute their respective thresholds. On the y -axis we plot the median, 50% and 95% of the epidemic threshold, computed over the instances.

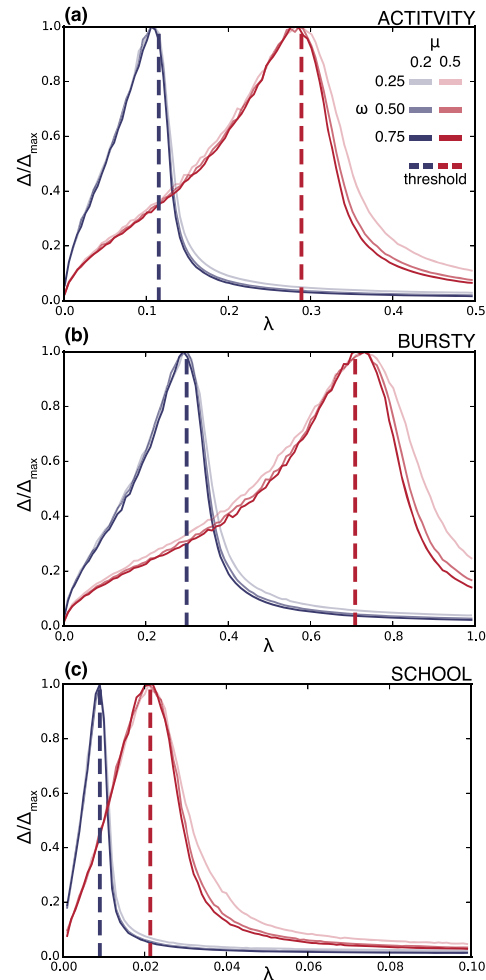


Fig. 5. SIRS model. Comparison between the epidemic threshold and variability. The variability Δ of the endemic prevalence i_T , normalized to its peak value (Δ/Δ_{\max}), is plotted against the transmissibility λ . We explore different values of μ and ω . Dashed vertical lines represent the threshold value for λ predicted by the infection propagator approach. (a) Shows the results for the ACTIVITY network, (b) for the BURSTY network, and (c) for the SCHOOL network.

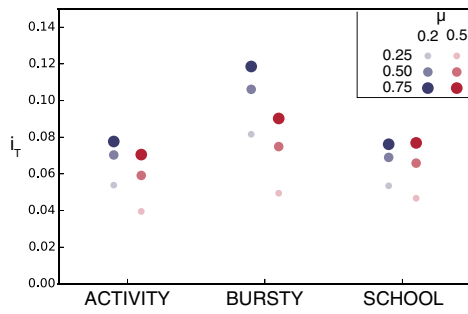


Fig. 6. SIRS model. Comparing the average prevalence above threshold, for different values of ω . For each network we choose a value of λ above threshold, so that variability is around $1/3$ of its peak, $\Delta/\Delta_{\max} \approx 1/3$. This value is clearly arbitrary, but it allows us to compare the different networks in similar epidemic situations. For ACTIVITY this corresponds to $\lambda = 0.14$ for $\mu = 0.2$, and $\lambda = 0.35$ for $\mu = 0.5$; for BURSTY $\lambda = 0.36$ and $\lambda = 0.86$, and for SCHOOL $\lambda = 0.01$ and $\lambda = 0.03$. We plot the average prevalence at equilibrium i_T for all three networks, and for all the explored values of μ and ω .

does not alter the accuracy of the computed predictions. Moreover, different immunity periods (i.e. different values of ω) lead to the same epidemic threshold on the temporal networks, as predicted by the infection propagator approach. The difference observed in the curves for different values of ω for λ well above the threshold is induced by the variation in the average endemic prevalence. Epidemics circulating on these systems and characterized by longer immunity periods ($\omega = 0.25$, light blue and light red in the plots of Fig. 5) display a larger variability due to the smaller average prevalence reached at equilibrium, as shown by Figure 6. Shorter immunity periods ($\omega = 0.75$) reach a larger endemic prevalence for a given value of the transmissibility above the epidemic threshold and therefore display a smaller variability Δ/Δ_{\max} .

5 Impact of time aggregation of the temporal network

In many cases, information on network dynamics can be coarse, with data reporting on the temporal evolution at a lower resolution scale than the one of the process itself. This means that all events occurring within the time interval of the considered resolution will be aggregated in a static single snapshot. An aggregated representation of a temporal network does not account for causal structures and temporal correlations that occur at time scales that are smaller than aggregation interval [59]. Since these structures can impact disease dynamics, it is crucial to assess how such coarser representation influences the description of epidemic processes [37,59,60]. Here, we study the influence of the aggregation schemes described in reference [60] on the epidemic threshold. HET scheme is a weighted aggregation of the snapshots, obtained by summing link weights: $\mathbf{W}_t, \mathbf{W}_{t+1} \mapsto \mathbf{W}_t + \mathbf{W}_{t+1}$. HOM is topologically equivalent to HET, having the exact same

set of links. Each link is given an equal weight corresponding to the average link weight of the weight distribution of the HET network aggregated over the same period. As a result, both schemes share the same average weight at every aggregation interval, but HET accounts for weight heterogeneity. We use the intrinsic transmissibility λ for comparison across different aggregation schemes and intervals, as it does not depend on weight. We consider the empirical dataset of the SCHOOL network as it provides a richer temporal and topological set of features with respect to synthetic models. Also, the study on aggregation aims at providing useful practical information for data collection purposes.

We consider the highest resolution network obtained from the SCHOOL data corresponding to $\Delta t_1 = 20$ s. Starting from this resolution, we aggregate snapshots recursively two by two, doubling the aggregation interval at each aggregation test. We consider the recovery rate m as an intrinsic property of the disease, thus not changing with aggregation. The probability of recovery after a time Δt is $m e^{-m \Delta t}$. Aggregation interval at the k th aggregation is $\Delta t_k = k \Delta t_1$. Hence, we compute the recovery probability at the k th aggregation $\mu[\Delta t_k]$ as the probability of recovering within an interval Δt_k , i.e., $\mu[\Delta t_k] = 1 - e^{-m \Delta t_1 k}$. We explore four different recovery rates, $m = 1.8, 9, 18, 90 \text{ h}^{-1}$, in order to explore different time scales of disease diffusion. High recovery rates mean short average infectious periods, and thus fast disease progression at node level. Conversely, low recovery rates induce long infectious periods, resulting in slower microscopical disease dynamics.

Threshold results on aggregated networks

We compare the epidemic threshold $\lambda_{\Delta t}$ computed after aggregating the network with a given aggregation time window Δt , to the one computed at the highest resolution Δt_1 (λ_1), using the ratio $\lambda_{\Delta t}/\lambda_1$ (Fig. 7). For each recovery rate, the results from two time aggregation schemes are shown.

Focusing on the HET aggregation scheme, the results of Figure 7a show that the prediction made on the aggregated SCHOOL network deteriorates with the increase of the time aggregation window Δt . As expected, the aggregation induces a loss of the temporal information making the aggregated network to perform poorly with respect to reproducing the behavior obtained in the original network. This is known for a series of indicators regarding the importance of individual nodes in the spread of an epidemic in the system [59], and we find that it also results in a biased estimation of the threshold condition for the epidemic propagation. The effect is more rapid and stronger for the fast disease (e.g. $m = 90 \text{ h}^{-1}$), in that it would have the possibility to experience the entire landscape of dynamical changes the network undergoes through, thus differentiating between the pattern obtained at the highest resolution and the aggregated one. On the other hand, for the slow disease (e.g. $m = 1.8 \text{ h}^{-1}$) we expect the epidemic process to be less sensitive to the network changes.

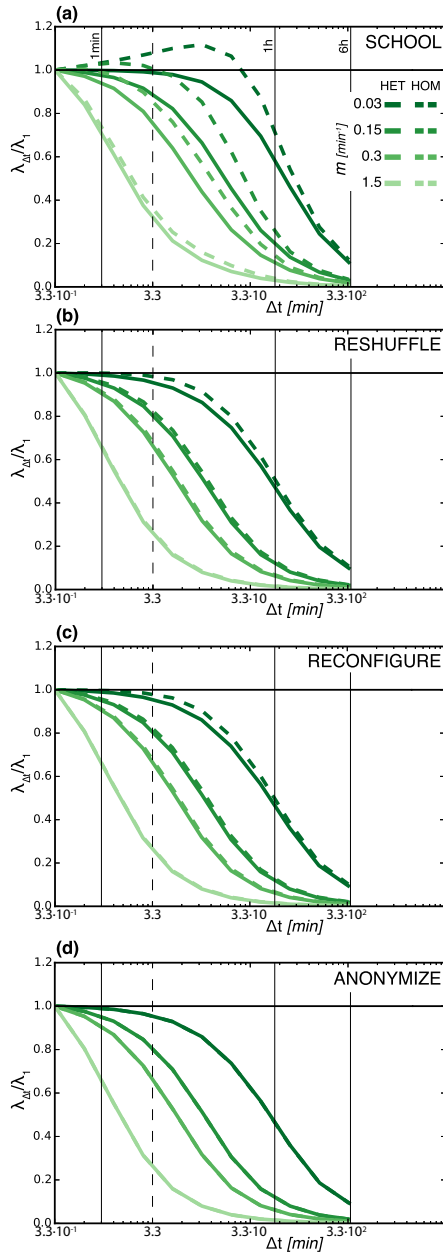


Fig. 7. Impact of aggregation on the epidemic threshold for the SCHOOL network and three reference models. The ratio between the threshold $\lambda_{\Delta t}$ computed on the aggregated network and the threshold λ_1 of the highest resolution temporal network is plotted as a function of the aggregation time interval Δt . Four different values of the recovery rate are explored, along with two aggregation schemes, HOM and HET. (a) Shows the results for SCHOOL, (b) for reference model RESHUFFLE, (c) for RECONFIGURE, and (d) for ANONYMIZE. Black vertical solid lines mark aggregation intervals of 1 min, 1 h and 6 h. Black vertical dashed line corresponds to 3 min 20 s.

The epidemic threshold computed on the aggregated network provides indeed a good estimate of the one corresponding to the highest resolution network up to a certain level of aggregation (e.g., $\Delta t \simeq 3$ min 20 s for the

$m = 1.8 \text{ h}^{-1}$ case), after which the accuracy is progressively lost. This is consistent with the numerical results of an SEIR (susceptible-exposed-infected-recovered) dynamics spreading on the network of contacts of conference attendees, showing that the spreading dynamics is well described by a static aggregated network if the heterogeneity of contact durations is taken into account as edge weights [60].

The underlying mechanism leading to the deterioration of the epidemic threshold estimate with increasing aggregation time window is the creation of novel transmission paths that would otherwise not exist, with the effect of destroying the causality of the sequence of interactions within Δt and of increasing the density of the links in the network [61,62]. All these effects tend to facilitate the spread of a disease, so that the resulting epidemic threshold is lower than the one computed on the original temporal network corresponding to $\Delta t = \Delta t_1 = 20$ s, as shown in Figure 7a.

If we focus on the HOM aggregation scheme, we observe that the epidemic threshold predicted for a given Δt is systematically higher than the one obtained in the HET scheme for the same Δt value (same color, dashed lines vs. continuous lines in Fig. 7a). The reason lies in the way weights are distributed over the links of the aggregated networks. While the HET scheme preserves the heterogeneity of the duration of the contacts, cumulating the duration of the interaction established by each pair of individuals, this information is lost in the HOM scheme as the total contact duration is homogeneously distributed among all contacts. Heterogeneity of the weights has a strong effect on the evolution of epidemics [45,63–71] and in many cases it has the effect of favoring the spread of diseases [21,72–74]. This results in a lower epidemic threshold than its homogeneous counterpart, for a given Δt . The faster the disease is, the smaller is the difference observed in the epidemic threshold obtained from the two aggregation schemes.

To better explore the various facets of the SCHOOL temporal network having an impact on the threshold condition, we also consider three reference models that systematically destroy some of the network properties. RESHUFFLE consists of a random reshuffling of snapshot time ordering. It preserves the aggregated network, and the static topological features of the snapshots. It breaks the temporal activity of the network, defined as the number of contacts in time. It breaks all temporal correlations among link activations, too. RECONFIGURE consists of a random reassignment of contact timestamps. Two contacts (i, j, t) , (k, l, s) are randomly selected, and their timestamp switched: (i, j, s) , (k, l, t) . It is equivalent to the DCW null model introduced in reference [18]. RECONFIGURE preserves the activity timeline and the aggregated network. It breaks snapshot topology and temporal correlations between link activations. Finally, ANONYMIZE reshuffles the identity of the nodes of each time snapshot, thus preserving activity timeline and static topology of each snapshot. It breaks all dynamic community structures and cliques (namely school classes).

Results for these reference models are shown in Figures 7b–7d. The behaviors obtained for RESHUFFLE and RECONFIGURE models are very similar. The difference between HET and HOM schemes is reduced for all recovery rates with respect to the results obtained on the original network, and it becomes negligible for faster diseases. The curves of Figures 7b and 7c show that the obtained result is independent of the activity timeline of the network (preserved by the RECONFIGURE reference model, but not by the RESHUFFLE one), and it is more likely related to specific time-evolving topological structures present in the SCHOOL network that are otherwise destroyed by both reference models. To test this hypothesis, we consider the ANONYMIZE reference model (Fig. 7d), where we destroy all two-points correlations and their time correlations, while preserving the overall temporal activity and the topology of each snapshot. As expected, the two schemes cannot be anymore distinguished following such reshuffling.

Results of Figure 7 show that in all reference models aggregation leads to an underestimation of the epidemic threshold, for both aggregation schemes considered. In the SCHOOL network, on the other hand, HOM aggregation is found to provide larger epidemic thresholds than the one obtained at the highest resolution, within a given aggregation interval and for slow diseases. To better understand this behavior observed solely on the empirical data that disappears with the three types of reshuffling considered in the reference models, we explore the role of time correlations and memory effects in the SCHOOL network. We consider the *social strategy* introduced in [75]. More in detail, we fix a time window of $\delta = 20$ snapshots (6 min 40 s), and define $k_{t,i}^{HOM}$ as the degree of node i in the network aggregated over the interval $[t - \delta, t]$. The degree assumes the same value in both networks as it is the number of incident connections. We also define $s_{t,i}^{HET}$ as the weighted degree of node i in HET network, i.e., the sum of the weights of its incident links [11]. We compute the social strategy of node i at time step t as $\gamma_{t,i} = k_{t,i}^{HOM} / s_{t,i}^{HET}$ (the same definition as in Ref. [75], except for a normalizing factor δ). Social strategy discriminates between memory-driven behavior ($\gamma \rightarrow 0$), where a node tends to make contacts always with the same nodes, and memoryless behavior ($\gamma \rightarrow 1$), where a node shows a more socially exploratory behavior. Figure 8a shows how social strategy evolves in time. We observe that its median behavior is quite stable in time around low values, except for several localized spikes. Most of these spikes roughly correspond to abrupt variations in the temporal activity of the network. These spikes result then from a reduction of the memory of the system, due to a varying number of overall contacts. Remarkably, however, the median social strategy returns to the value it had before the spike quickly after each of these events, indicating that the interaction dynamics does not qualitatively change, but the sets of interacting individuals do change over time. The only exception occurs between around 13 : 30 and 14 : 30, when social strategy is significantly higher than average, but still lower than its delimiting peaks. These spikes

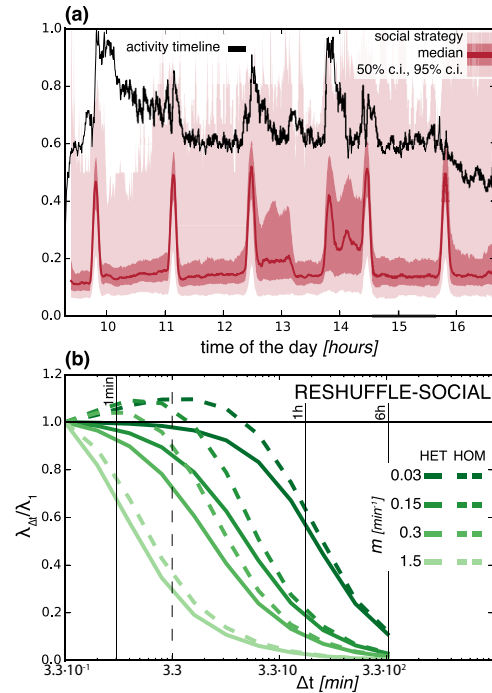


Fig. 8. Interplay between SCHOOL dynamics and aggregation. In (a) we show the temporal activity of SCHOOL, as the normalized number of contacts (black line) in each snapshot of the fully temporal network. On the x -axis we indicate the time of the day, in hours. Red solid line represents the median value of network’s social strategy, computed with a sliding window of 20 snapshots (equivalent to 400 s). Social strategy plotted at time t is computed over the interval $(t - 20, t]$. Red areas show 50% (darker) and 95% (lighter) confidence interval for social strategy. In (b) we plot, for reference model RESHUFFLE-SOCIAL, the ratio between the threshold computed on the aggregated network $\lambda_{\Delta t}$, and the threshold of the full temporal network λ_1 , as a function of the aggregation time interval Δt . Four different values of the recovery rate are explored, along with the two aggregation schemes HOM and HET. Black vertical solid lines mark aggregation intervals of 1 min, 1 h and 6 h. Black vertical dashed line corresponds to 3 min 20 s.

naturally induce a temporal slicing of the network, in a way that likely corresponds to the rhythm of school activities. We call γ -slice each time interval between two consecutive spikes.

The degree of memory contained in the system, and measured through the tendency of each node to keep establishing contacts with the same individuals over time, is destroyed in all reference models under study, even those that preserve the activity timeline of Figure 8a. To understand whether if and to what extent stages in the time evolution of the social strategy are responsible for the behavior observed in the SCHOOL network, we design a fourth reference model, RESHUFFLE-SOCIAL, where we randomize the snapshot order, as in RESHUFFLE, but we allow reshuffling only within each γ -slice. Figure 8b shows that RESHUFFLE-SOCIAL displays the same behavior as the SCHOOL network, unlike RESHUFFLE, with an overestimation of the value of the epidemic threshold by the HOM

scheme for small enough aggregation intervals and slow diseases. The aggregation of snapshots where individuals show a rather large memory in the way they establish links (i.e. small γ) leads to marked weight-topology correlations, likely being part of robust temporal communities of highly interacting nodes emerging from school daily activities. Such correlations were already found to play an important role in the slowing down of epidemics once large-scale propagation occurs in the system [18]. In our case, we find that preserving the heterogeneity of weights of such correlations (as in the HET scheme) can provide a good approximation of the epidemic threshold for small interval aggregation and for slow diseases. In addition, such approximation is better than the one provided by homogenizing weights across all links in the system (as in the HOM scheme), given that the latter destroys weight-topology correlations leading to a network that is more resilient to the epidemic spread [21]. This effect vanishes for increasing time aggregating windows and it completely disappears for all aggregating intervals once these correlations are destroyed by the reshuffling of nodes (see the ANONYMIZE reference model in Fig. 7).

6 Conclusions

We have considered the infection propagator approach to compute the epidemic threshold for an arbitrary time-varying network. Starting from a SIS dynamics on a weighted directed temporal network, we have considered more complicated compartmental models and addressed timescale issues relevant for the study of temporal networks. The overall aim was to introduce the infection propagator approach for more realistic infection dynamics and to study the effect of time aggregation of the network of contacts on the computation of its threshold. Our findings indicate that the approach provides reliable and accurate predictions of the epidemic threshold also in presence of immunity stages and loss of immunity transitions in the disease natural history. In addition, for slow diseases, the time aggregation scheme preserving the cumulative heterogeneous duration of contacts between two nodes is shown to provide a quite accurate estimation of the epidemic threshold of the corresponding high-resolution network up to a certain aggregation level. For faster diseases, time aggregation strongly alters the accuracy of the estimation. The presence of weight-topology correlations is the main feature of the SCHOOL network leading to biased estimations. These findings provide important information to study the vulnerability of systems in real settings and to assess possible biases induced by the consideration of time-aggregated contact data.

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References

1. R.M. Anderson, R.M. May, *Infectious Diseases of Humans Dynamics and Control* (Oxford University Press, 1992)
2. M.J. Keeling, P. Rohani, *Modeling Infectious Diseases in Humans and Animals* (Princeton University Press, 2008)
3. R. Pastor-Satorras, A. Vespignani, *Phys. Rev. Lett.* **86**, 3200 (2001)
4. Y. Wang, D. Chakrabarti, C. Wang, C. Faloutsos, Epidemic spreading in real networks: an eigenvalue viewpoint, in *Proceedings of 22nd International Symposium on Reliable Distributed Systems, 2003* (2003), pp. 25–34
5. S. Gómez, A. Arenas, J. Borge-Holthoefer, S. Meloni, Y. Moreno, *Europhys. Lett.* **89**, 38009 (2010)
6. R. Cohen, K. Erez, D. ben Avraham, S. Havlin, *Phys. Rev. Lett.* **85**, 4626 (2000)
7. M.E.J. Newman, *Phys. Rev. E* **66**, 016128 (2002)
8. M. Boguñá, C. Castellano, R. Pastor-Satorras, *Phys. Rev. Lett.* **111**, 068701 (2013)
9. C. Castellano, R. Pastor-Satorras, *Phys. Rev. Lett.* **105**, 218701 (2010)
10. A.V. Goltsev, S.N. Dorogovtsev, J.G. Oliveira, J.F.F. Mendes, *Phys. Rev. Lett.* **109**, 128702 (2012)
11. A. Barrat, M. Barthélemy, A. Vespignani, *Dynamical Processes on Complex Networks* (Cambridge University Press, 2008)
12. S. Bansal, J. Read, B. Pourbohloul, L.A. Meyers, *J. Biol. Dyn.* **4**, 478 (2010)
13. L. Isella, J. Stehlé, A. Barrat, C. Cattuto, J.F. Pinton, W. Van den Broeck, *J. Theor. Biol.* **271**, 166 (2011)
14. G. Miritello, E. Moro, R. Lara, *Phys. Rev. E* **83**, 045102 (2011)
15. L.E.C. Rocha, F. Liljeros, P. Holme, *Proc. Natl. Acad. Sci. USA* **107**, 5706 (2010)
16. P. Bajardi, A. Barrat, F. Natale, L. Savini, V. Colizza, *PloS One* **6**, e19869 (2011)
17. C.T. Butts, *Science* **325**, 414 (2009)
18. M. Karsai, M. Kivela, R.K. Pan, K. Kaski, J. Kertész, A.L. Barabási, J. Saramäki, *Phys. Rev. E* **83**, 025102 (2011)
19. P. Holme, J. Saramäki, *Phys. Rep.* **519**, 97 (2012)
20. L.E.C. Rocha, V.D. Blondel, *PLoS Comput. Biol.* **9**, e1002974 (2013)
21. L. Ferreri, P. Bajardi, M. Giacobini, S. Perazzo, E. Venturino, *Phys. Rev. E* **90**, 012812 (2014)
22. J.L. Iribarren, E. Moro, *Phys. Rev. Lett.* **103**, 038702 (2009)
23. A. Vazquez, B. Rácz, A. Lukács, A.L. Barabási, *Phys. Rev. Lett.* **98**, 158702 (2007)
24. K.T. Eames, M.J. Keeling, *Math. Biosci.* **189**, 115 (2004)
25. N. Perra, B. Gonçalves, R. Pastor-Satorras, A. Vespignani, *Sci. Rep.* **2**, 469 (2012)
26. E. Volz, L.A. Meyers, *J. Roy. Soc. Int.* **6**, 233 (2009)
27. T. Gross, C.J.D. D’Lima, B. Blasius, *Phys. Rev. Lett.* **96**, 208701 (2006)
28. M. Taylor, T.J. Taylor, I.Z. Kiss, *Phys. Rev. E* **85**, 016103 (2012)
29. Z. Zhao, J.P. Calderón, C. Xu, G. Zhao, D. Fenn, D. Sornette, R. Crane, P.M. Hui, N.F. Johnson, *Phys. Rev. E* **81**, 56107 (2010)
30. E. Valdano, L. Ferreri, C. Poletto, V. Colizza, *Phys. Rev. X* **5**, 21005 (2015)

31. M. De Domenico, A. Solé-Ribalta, E. Cozzo, M. Kivela, Y. Moreno, M.A. Porter, S. Gómez, A. Arenas, *Phys. Rev. X* **3**, 041022 (2013)
32. M. Kivela, A. Arenas, M. Barthelemy, J.P. Gleeson, Y. Moreno, M.A. Porter, *J. Complex Networks* **2**, 203 (2014)
33. P. Vanhems, A. Barrat, C. Cattuto, J.F. Pinton, N. Khanafer, C. Régis, B. Kim, B. Comte, N. Voirin, *PLoS One* **8**, e73970 (2013)
34. M. Korschake, H.H.K. Lentz, F.J. Conraths, P. Hövel, T. Selhorst, *PLoS One* **8**, e55223 (2013)
35. J. Fournet, A. Barrat, *PLoS One* **9**, e107878 (2014)
36. T. Obadia, R. Silhol, L. Opatowski, L. Temime, J. Legrand, A.C.M. Thiébaud, J.L. Herrmann, E. Fleury, D. Guillemot, P.Y. Boëlle, *PLoS Comput. Biol.* **11**, e1004170 (2015)
37. B. Ribeiro, N. Perra, A. Baronchelli, *Sci. Rep.* **3**, 3006 (2013)
38. P. Bajardi, A. Barrat, L. Savini, V. Colizza, *J. Roy. Soc. Int.* **9**, 2814 (2012)
39. A. Machens, F. Gesualdo, C. Rizzo, A. Tozzi, A. Barrat, C. Cattuto, *BMC Infectious Diseases* **13**, 185 (2013)
40. J.P. Gleeson, S. Melnik, J.A. Ward, M.A. Porter, P.J. Mucha, *Phys. Rev. E* **85**, 026106 (2012)
41. P. Van Mieghem, *Computing* **93**, 147 (2011)
42. P. Van Mieghem, R. van de Bovenkamp, *Phys. Rev. Lett.* **110**, 108701 (2013)
43. C. Li, H. Wang, P. Van Mieghem, *Phys. Rev. E* **88**, 062802 (2013)
44. K. Wehmuth, A. Ziviani, E. Fleury, [arXiv:1402.3488](https://arxiv.org/abs/1402.3488) (2014)
45. P.J. Mucha, T. Richardson, K. Macon, M.A. Porter, J.P. Onnela, *Science* **329**, 277 (2010)
46. I. Scholtes, N. Wider, R. Pfitzner, A. Garas, C.J. Tessone, F. Schweitzer, *Nat. Commun.* **5**, 5024 (2014)
47. E. Cozzo, R.A. Baños, S. Meloni, Y. Moreno, *Phys. Rev. E* **88**, 050801 (2013)
48. H. Wang, Q. Li, G. D'Agostino, S. Havlin, H.E. Stanley, P. Van Mieghem, *Phys. Rev. E* **88**, 022801 (2013)
49. C. Granell, S. Gómez, A. Arenas, *Phys. Rev. Lett.* **111**, 128701 (2013)
50. S. Elaydi, *An Introduction to Difference Equations*, 3rd edn. (Springer, New York, 2005)
51. P.D. Powell, [arXiv:1112.4379](https://arxiv.org/abs/1112.4379) (2011)
52. H.K. Lentz, T. Selhorst, I.M. Sokolov, *Phys. Rev. Lett.* **110**, 118701 (2013)
53. P. Van Mieghem, F. Sahneh, C. Scoglio, An upper bound for the epidemic threshold in exact Markovian SIR and SIS epidemics on networks, in *IEEE 53rd Annual Conference on Decision and Control (CDC), 2014* (2014), pp. 6228–6233
54. P. Shu, W. Wang, M. Tang, Y. Do, *Chaos* **25**, 063104 (2015)
55. M. Salathé, M. Kazandjeva, J.W. Lee, P. Levis, M.W. Feldman, J.H. Jones, *Proc. Natl. Acad. Sci USA* **107**, 22020 (2010)
56. S.C. Ferreira, C. Castellano, R. Pastor-Satorras, *Phys. Rev. E* **86**, 041125 (2012)
57. P. Crépey, F.P. Alvarez, M. Barthélemy, *Phys. Rev. E* **73**, 046131 (2006)
58. S.C. Ferreira, R.S. Ferreira, C. Castellano, R. Pastor-Satorras, *Phys. Rev. E* **84**, 066102 (2011)
59. P. Holme, *PLoS Comput Biol* **9**, e1003142 (2013)
60. J. Stehlé, N. Voirin, A. Barrat, C. Cattuto, V. Colizza, L. Isella et al., *BMC Medicine* **9**, 87 (2011)
61. P. Holme, *Phys. Rev. E* **71**, 046119 (2005)
62. R.K. Pan, J. Saramaki, *Phys. Rev. E* **84**, 016105 (2011)
63. M.S. Granovetter, *Am. J. Sociol.* **78**, 1360 (1973)
64. Y. Gang, Z. Tao, W. Jie, F. Zhong-Qian, W. Bing-Hong, *Chinese Phys. Lett.* **22**, 510 (2005)
65. J.P. Onnela, J. Saramäki, J. Hyvönen, G. Szabó, D. Lazer, K. Kaski, J. Kertész, A.L. Barabási, *Proc. Natl. Acad. Sci.* **104**, 7332 (2007)
66. R. Lambiotte, J.C. Delvenne, M. Barahona, [arXiv:0812.1770](https://arxiv.org/abs/0812.1770) (2009)
67. R. Toivonen, X. Castelló, V.M. Eguíluz, J. Saramäki, K. Kaski, M. San Miguel, *Phys. Rev. E* **79**, 016109 (2009)
68. K.T. Eames, J.M. Read, W.J. Edmunds, *Epidemics* **1**, 70 (2009)
69. M. Salathé, J.H. Jones, *PLoS Comput. Biol.* **6**, e1000736 (2010)
70. Z. Yang, T. Zhou, *Phys. Rev. E* **85**, 056106 (2012)
71. C. Kamp, M. Moslonka-Lefebvre, S. Alizon, *PLoS Comput. Biol.* **9**, e1003352 (2013)
72. M. Deijfen, *Math. Biosc.* **232**, 57 (2011)
73. T. Britton, M. Deijfen, F. Liljeros, *J. Stat. Phys.* **145**, 1368 (2011)
74. T. Britton, D. Lindenstrand, *Math. Biosc.* **240**, 124 (2012)
75. G. Miritello, R. Lara, M. Cebrian, E. Moro, *Sci. Rep.* **3**, 1950 (2013)